

ANALYTICITY OF GREEN'S FUNCTIONS OF DILUTE QUANTUM GASES

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In this note we point out that Ginibre's results on the reduced density matrices of quantum gases¹⁾ have immediate implications for the existence and analyticity of Green's functions. If H_Λ is the Hamiltonian in the bounded region Λ , we define Green's functions by

$$G_\Lambda(x_1, \dots, x_m; \zeta_1, \dots, \zeta_m) = Z^{-1} \text{Tr}(A_1(x_1) e^{-(\zeta_2 - \zeta_1)H_\Lambda} A_2(x_2) \dots e^{-(\zeta_m - \zeta_{m-1})H_\Lambda} A_m(x_m) e^{-(\beta + \zeta_1 - \zeta_m)H_\Lambda})$$

where $Z = \text{Tr} e^{-\beta H_\Lambda}$ and $A_k(x_k) = a^*(x'_{k1}) \dots a^*(x'_{kp(k)}) a(x''_{k1}) \dots a(x''_{kq(k)})$.

Let $\varphi_k \in L^2((\mathbb{R}^V)^{p(k)+q(k)})$ for Fermi statistics, or $\varphi_k = \varphi'_k \varphi''_k$ with $\varphi'_k \in L^2(\mathbb{R}^{Vp(k)})$, $\varphi''_k \in L^2(\mathbb{R}^{Vq(k)})$ for Bose statistics; we write

$$G_\Lambda^\varphi(\zeta_1, \dots, \zeta_m) = \int_{\Lambda^{p(1)+q(1)}} dx_1 \dots \int_{\Lambda^{p(m)+q(m)}} dx_m \varphi_1(x_1) \dots \varphi_m(x_m) G_\Lambda(x_1, \dots, x_m; \zeta_1, \dots, \zeta_m)$$

In the case of a system of particles interacting through a suitable pair potential Φ and for small activity the operator $e^{-\lambda H_\Lambda}$, with $\lambda > 0$, may be defined in terms of Wiener integrals and is of trace class. The operators $e^{-\lambda H_\Lambda} A_k(x_k) e^{-\lambda H_\Lambda}$ can also be expressed in terms of Wiener integrals and are of trace class.

When λ is complex and $\text{Re } \lambda > 0$, $e^{-\lambda H_\Lambda}$ is defined and analytic, therefore G_Λ is an analytic function of the complex variables $\zeta_k = \beta_k - it_k$ in the domain

$$\mathfrak{D} = \{(\zeta_1, \dots, \zeta_m) : \beta_1 < \dots < \beta_m < \beta_1 + \beta\}$$

If $t_1 = \dots = t_m$, and $\beta_1 < \dots < \beta_m < \beta_1 + \beta$, G_Λ can be expressed in terms of Wiener integrals and it follows from Ginibre's analysis²⁾ that when $\Lambda \rightarrow \infty$ (e.g. Λ is a sphere centered at the origin and with radius tending to infinity),

$$G_\Lambda(x_1, \dots, x_m; \beta_1, \dots, \beta_m) \rightarrow G(x_1, \dots, x_m; \beta_1, \dots, \beta_m) \quad (1)$$

uniformly on compacts with respect to x_1, \dots, x_m , and

$$\begin{aligned} G_\Lambda^\varphi(\beta_1, \dots, \beta_m) &\rightarrow G^\varphi(\beta_1, \dots, \beta_m) \\ &= \int dx_1 \dots dx_m \varphi_1(x_1) \dots \varphi_m(x_m) G(x_1, \dots, x_m; \beta_1, \dots, \beta_m) \end{aligned} \quad (2)$$

1. Proposition. There exists a function $G(x_1, \dots, x_m; \zeta_1, \dots, \zeta_m)$ analytic with respect to $(\zeta_1, \dots, \zeta_m) \in \mathcal{D}$ and such that

$$\lim_{\Lambda \rightarrow \infty} G_\Lambda(x_1, \dots, x_m; \zeta_1, \dots, \zeta_m) = G(x_1, \dots, x_m; \zeta_1, \dots, \zeta_m)$$

uniformly on compacts with respect to $x_1, \dots, x_m, \zeta_1, \dots, \zeta_m$.

Furthermore, if $(\zeta_1, \dots, \zeta_m) \in \mathcal{D}$,

$$\begin{aligned} \lim_{\Lambda \rightarrow \infty} G_\Lambda^\varphi(\zeta_1, \dots, \zeta_m) &= G^\varphi(\zeta_1, \dots, \zeta_m) \\ &= \int dx_1 \dots dx_m \varphi_1(x_1) \dots \varphi_m(x_m) G(x_1, \dots, x_m; \zeta_1, \dots, \zeta_m) \end{aligned}$$

We notice first that \mathcal{D} is the union (over $n > 0$) of the sets

$$\mathcal{H}_n = \{(\zeta_1, \dots, \zeta_m) \in \mathcal{D} : \beta_2 - \beta_1 \geq \frac{\beta}{2n}, \dots, \beta_m - \beta_{m-1} \geq \frac{\beta}{2n}, \beta + \beta_1 - \beta_m \geq \frac{\beta}{2n}\}$$

If $(\xi_1, \dots, \xi_m) \in \mathcal{H}_n$ we may express G_Λ in terms of operators

$$e^{(it_k - \frac{1}{4n} \beta)H_\Lambda} A_k(x_k) e^{-(it_k + \frac{1}{4n} \beta)H_\Lambda}, \text{ and } e^{-\lambda H_\Lambda} \text{ with } 0 < \lambda < \beta.$$

Using Hölder's inequality³⁾ we find an upper bound for $|G_\Lambda|$ in terms of the expressions

$$Z^{-1} \text{Tr} \left[\left(e^{-\frac{1}{4n} \beta H_\Lambda} A_k(x_k) e^{-\frac{1}{2n} \beta H_\Lambda} A_k(x_k) e^{-\frac{1}{4n} \beta H_\Lambda} \right)^n \right]$$

which are known by (1) to have a limit when $\Lambda \rightarrow \infty$. We may thus assume that G_Λ is bounded on each \mathcal{H}_n uniformly with respect to Λ and x_1, \dots, x_m in a compact, and the convergence of G_Λ on the real points of \mathcal{D} implies its uniform convergence on the compacts of \mathcal{D} . The uniformity of the convergence with respect to x_1, \dots, x_m on compacts follows from the uniformity of (1).

The proof of the convergence of G_Λ^φ proceeds like the proof of the convergence of G_Λ and shows in particular that the limit of G_Λ^φ is a bounded multilinear functional of $\varphi_1, \dots, \varphi_m$ (Fermi) or $\varphi_1', \dots, \varphi_m''$ (Bose) on the product of the relevant L^2 spaces, identification of the limit follows from (2), taking $\varphi_1, \dots, \varphi_m$ with compact supports.

2. Proposition. Let $m=2$ and let the pair potential
 $\Phi \in L^1(\mathbb{R}^\nu) \cap L^2(\mathbb{R}^\nu)$. Then G^φ extends to a bounded continuous function
on \mathcal{D} such that

$$\lim_{\Lambda \rightarrow \infty} G_\Lambda^\varphi(\xi_1, \xi_2) = G^\varphi(\xi_1, \xi_2)$$

uniformly on the compacts of the closure $\overline{\mathfrak{D}}$ of \mathfrak{D} ,

The operators $A_k(\varphi_k) e^{-\lambda H_{\wedge}}$ are of trace class [consider $e^{-\lambda H_{\wedge}} A_k(\varphi_k)^* A_k(\varphi_k) e^{-\lambda H_{\wedge}}$] and, if $\beta_2 = \beta_1$ or $\beta_2 = \beta_1 + \beta$, we have

$$|G_{\wedge}^{\varphi}(\xi_1, \xi_2)| \leq |G_{1\wedge} G_{2\wedge}|^{1/2} \quad (3)$$

$$G_{k\wedge} = Z^{-1} \text{Tr}[(A_k(\varphi_k)^* A_k(\varphi_k) + A_k(\varphi_k) A_k(\varphi_k)^*) e^{-\beta H_{\wedge}}] \quad (4)$$

We assume first Bose statistics. The reduced density matrices are integral kernels of bounded operators in L^2 . When (\wedge_n) tends to infinity these bounded operators form a bounded sequence converging in the strong operator topology¹⁾. Therefore there exists $C > 0$ such that

$$G_{k\wedge_n} \leq C \int dx_k |\varphi_k(x_k)|^2 \quad (5)$$

for all n . Since G_{\wedge}^{φ} is analytic and bounded, (3) holds for all (ξ_1, \dots, ξ_2) in $\overline{\mathfrak{D}}$ and, using (5), this gives

$$|G_{\wedge_n}^{\varphi}(\xi_1, \xi_2)| \leq C \|\varphi_1\|_2 \|\varphi_2\|_2 \quad (6)$$

for all n and $(\xi_1, \xi_2) \in \overline{\mathfrak{D}}$. In the case of Fermi statistics, (6) holds again (with $C=1$) because

$$\|A_k(\varphi_k)\| \leq \|\varphi_k\|_2$$

In view of (6) it suffices to prove the proposition when $A_k(\varphi_k)$ is of the form

$$A_k(\varphi_k) = a^*(\psi_{k1}') \dots a^*(\psi_{kp(k)}') a(\psi_{k1}'') \dots a(\psi_{kq(k)}'')$$

where $\psi_{k1}', \dots, \psi_{kq(k)}''$ are of class C^2 with compact support.

We have

$$\frac{d}{d\zeta_2} G_{\wedge}^{\varphi} = Z^{-1} \text{Tr} \{ A_1(\varphi_1) e^{-(\zeta_2 - \zeta_1) H_{\wedge}} [A_2(\varphi_2), H_{\wedge}] e^{-(\beta + \zeta_1 - \zeta_2) H_{\wedge}} \}$$

and therefore

$$\left| \frac{d}{d\zeta_2} G_{\wedge}^{\varphi} \right| \leq [G_{1\wedge} G_{2\wedge}']^{1/2}$$

where $G_{2\wedge}'$ is given by (4) with $A_k(\varphi_k)$ replaced by $[A_2(\varphi_2), H_{\wedge}]$. Since $\psi_{21}', \dots, \psi_{2q(2)}''$ are of class C^2 with compact support, the commutator of $A_2(\varphi_2)$ with the kinetic energy part of H_{\wedge} is again of the form $A(\varphi)$. In view of this $G_{2\wedge}'$ is a sum of integrals of reduced density matrices $G_{\wedge}'(x_1, \dots, x_r)$ multiplied by continuous functions $\psi(x_i)$ with compact support, and the pair potential $\Phi(x_k - x_j)$. The pair potential appears as factor 0, 1 or 2 times; if $\Phi(x_k - x_j)$ appears there also appears a factor $\psi(x_j)$ or $\psi(x_k)$; for each variable x_i in $G_{\wedge}'(x_1, \dots, x_r)$ which does not appear in a factor $\Phi(x_j - x_i)$ there is a factor $\psi(x_i)$. Using the condition $\Phi \in L^1(\mathbb{R}^V) \cap L^2(\mathbb{R}^V)$ and the fact that the reduced density matrices G_{\wedge}' are bounded functions, uniformly in $\wedge^{(1)}$ we obtain a bound on $G_{2\wedge}'$ which is independent of \wedge .

Therefore $\left| \frac{d}{d\zeta_2} G_{\wedge}^{\varphi} \right| = \left| \frac{d}{d\zeta_1} G_{\wedge}^{\varphi} \right|$ is bounded on \mathfrak{D} uniformly in \wedge . The convergence of G_{\wedge}^{φ} in \mathfrak{D} implies then its uniform convergence on the compacts of $\overline{\mathfrak{D}}$.

3. Remark. Let $m=3$, $\Phi \in L^1(\mathbb{R}^V) \cap L^2(\mathbb{R}^V)$. In the case of Fermi statistics, G^{φ} extends to a bounded continuous function on $\overline{\mathfrak{D}}$ such that

$$\lim_{\wedge \rightarrow \infty} G_{\wedge}^{\varphi}(\zeta_1, \zeta_2, \zeta_3) = G^{\varphi}(\zeta_1, \zeta_2, \zeta_3)$$

uniformly on the compacts of the closure $\overline{\mathfrak{D}}$ of \mathfrak{D} .

To estimate $\left| \frac{d}{d\zeta_1} G_{\wedge}^{\varphi} \right|$ it suffices to consider the expression

$$Z^{-1} \text{Tr} \{ e^{(-\beta+it)H_{\wedge}} [A_i, H_{\wedge}] e^{it'H_{\wedge}} A_j e^{it''H_{\wedge}} A_k \}$$

and similar ones where $[A_i, H_{\wedge}]$ and A_j, A_k are circularly permuted. If $[A_i, H_{\wedge}]$ occupies the middle position we rewrite the expression in terms of $[A_j, H_{\wedge}], [A_k, H_{\wedge}]$. The rest of the argument goes as for $m=2$ (using the boundedness of A_j).

4. Proposition. In the case of Fermi statistics, introduce the operators

$$A_k(\varphi_k, f_k) = \int dt f(t) e^{it H_{\wedge}} A_k(\varphi_k) e^{-it H_{\wedge}}$$

where $f_k \in L^1(\mathbb{R})$, then the limit

$$\lim_{\Lambda \rightarrow \infty} Z^{-1} \text{Tr}(A_1(\varphi_1, f_1) \dots A_m(\varphi_m, f_m) e^{-\beta H_\Lambda})$$

exists.

It is sufficient to prove this for f_k of class C^1 with compact support. We construct Green's functions with the operators $A_k(\varphi_k, f_k)$ instead of $A_k(\varphi_k)$ and use the fact that the derivatives of these functions are bounded in \mathcal{D} uniformly with respect to Λ .

5. Remark. Proposition 4 has obvious implications for the description of time evolution of a dilute Fermi gas. It does not, however, exhibit this time evolution as a group of automorphisms of the C^* -algebra of the anticommutation relations. Streater⁴⁾ and Hepp⁵⁾ have shown that such a group of automorphisms exists for some non local interactions.

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Footnotes.

1) See J. Ginibre. J. Math. Phys. 6, 238-251 (1965); 6, 252-262 (1965); 6; 1432-1446 (1965); J. Ginibre p. 148 in Statistical Mechanics (proceedings of the I.U.P.A.P. meeting, Copenhagen, 1966) edited by T. Bak, Benjamin, New York, 1967.

2) This result is contained in C. Gruber's thesis [Princeton, 1968, unpublished], see also J. Ginibre and C. Gruber, Commun. Math. Phys. 11, 198-213 (1969).

3) See N. Dunford and J. Schwartz, Linear Operators, Interscience, New York, 1963, Lemma XI. 9-20, p 1105.

4) R.F. Streater. Commun. Math. Phys. 7, 93-98 (1968)

5) K. Hepp. Unpublished.